

# Characterization of new ferromagnetic Fe-Co-Zn-Ga alloys by ab initio investigations

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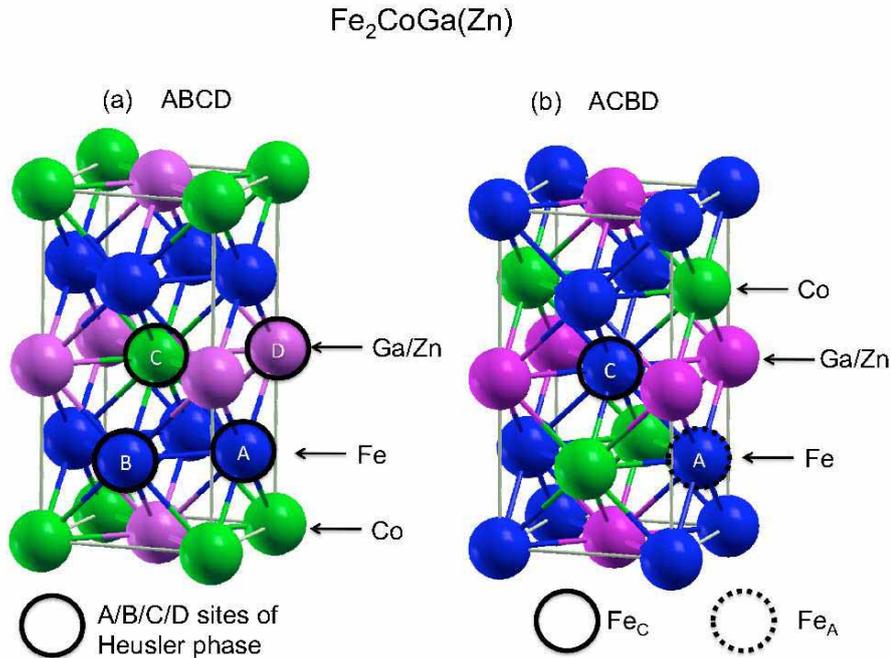
**Abstract.** We have calculated structural energy differences, magnetic interaction constants and mixing energies of Fe-Co-Ga(Zn) alloys in order to gain information about their suitability for ferromagnetic shape memory devices. We considered the classical Heusler structure as well as an ordering type in which the Co and one Fe sublattice are interchanged. The results of our density functional theory investigations suggest high Curie and martensitic phase transition temperatures for the Ga-based as well as for the Zn-based systems. In case of hypothetical Fe<sub>2</sub>CoZn alloys, the classically ordered Heusler structure is energetically preferred. In Ga based alloys, the ordering type with Co partially on Fe sites appears to be more stable. We propose that a systematic variation of composition by successive addition of Zn and Co to Fe-Co-Ga may result in a promising new ferromagnetic shape memory alloy of type Fe<sub>2-x</sub>Co<sub>1+x</sub>Ga<sub>1-y</sub>Zn<sub>y</sub>.

## 1 Introduction

Ferromagnetic shape memory alloys (FSMAs) have attracted considerable attention due to their potential use for new actuator materials. For prototype Ni<sub>2</sub>MnGa Heusler compound, large magnetically induced strains have been reported in the martensitic phase. However, the martensitic phase transition takes place at too low temperatures for many practical applications. Furthermore, Ni-Mn-Ga alloys are disadvantageous because of their poor room temperature ductility or high costs of elements [1–3]. In the literature, various other systems including Ni-Mn-(Al, Sn, Sb, In) [4, 5], Co-Ni-(Al, Ga) [6–13], Ni-Fe-(Al, Ga) [14] and Cu-Mn-(Al, Ga) [15, 16] have been proposed as candidates for new FSMAs. Of special interest are Heusler systems based on Co, Fe, and Zn: For the Co-Ni-(Al, Ga) and Ni-Ga-Fe-Co alloys, the martensitic start temperature  $T_S$  decreases with increasing Co content while Curie temperature and spontaneous magnetization increase [17, 18]. For Ni-Mn-Ga alloys, the martensitic transformation temperature  $T_M$  can be increased by changing the stoichiometry. Recently developed ferromagnetic Ni-Fe-Ga alloys show good plasticity of the high-temperature phase in comparison to Ni<sub>2</sub>MnGa and thus may be promising. In Fe-Co-Zn, the magnetic phase transition induces various interesting anomalies like a miscibility gap between the ferromagnetic and the paramagnetic bcc  $\alpha$ -phase [19]. The nonmagnetic Cu-Al-Zn Heusler alloys are important SMA [20, 21] while Fe-Ga is a magnetostrictive material [22, 23]. The compound Co<sub>2</sub>FeGa exhibits the classically ordered Heusler structure as early experimental studies report [24], while in Fe-Co-Ga, the Co atoms prefer to occupy parts of the Fe sublattices (e.g., the B site, see Fig. 1). Thus, one has to differentiate two types of Fe atoms: Fe atoms which remain on the Fe sublattice (e.g., the A site, here denoted

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**Fig. 1.** Left: Tetragonal, body-centered unit cell in the  $\text{Fe}_2\text{CoGa}(\text{Zn})$  Heusler structure, named ABCD in the order of appearance in the chemical formula. Right: Heusler unit cell with Co atoms (green spheres) on one Fe (blue spheres) sublattice, marked as ACBD. Here, half of the Fe atoms (B sites) occupy the Co sites (C sites and thus are named as  $\text{Fe}_C$ ) while the remaining half of the Fe atoms stay at their sites, and are consequently denoted by  $\text{Fe}_A$ . Ga(Zn) atoms are shown as magenta spheres.

by  $\text{Fe}_A$ ) and Fe atoms which occupy Co sites (the C site, denoted by  $\text{Fe}_C$ ) (see Fig. 1). At 25% Co, ordering type ACBD, with the elements placed on the sites in the respective order of their appearance in the formula unit (cf. Fig. 1), is achieved and with further increasing Co content the Fe sublattices becomes entirely occupied by Co atoms, so that we obtain the order ABCD for  $\text{Co}_2\text{FeGa}$ . In general, transition metal atoms to the right of Fe in the periodic table preferentially go to the Fe sublattice (e.g. B site) while those to the left of Fe go to the Co sites (e.g. C site) [24]. A considerable Curie temperature of the order of 1160 K is found for  $\text{Fe}_2\text{CoGa}$ .

These findings motivate a theoretical investigation of Fe-Co-Ga(Zn) alloys in the classical Heusler structure (ABCD) and in the ordering type ACBD (see Fig. 1). Quaternary alloys of type Fe-Ni-Co-Ga-Zn may be promising in order to increase the range of potential application of FSMA and for an appropriate choice of composition, high transition temperatures and an improvement of mechanical properties may be obtained.

## 2 Computational details

Our investigations were carried out in the framework of density functional theory using the plane wave code VASP [25]. The electron-ionic core interaction is accounted for by applying the PAW method [26] and the GGA exchange-correlation potential of Perdew, Burke, and Ernzerhoff [27]. The integration over the Brillouin zone was carried out using the tetrahedron method with Blöchl corrections and  $\Gamma$ -centered k-point grids of  $13 \times 13 \times 13$  (total energy) and

**Table 1.** Calculated results of tetragonally distorted Fe-Co-Ga(Zn),  $c/a \neq 1$ , optimum lattice constant,  $a$ , magnetic moment,  $M$ , mixing energy,  $E_{\text{mix}}$ , energy difference between the energy minimum at  $c/a < 1$  and the energy minimum at  $c/a > 1$ ,  $\Delta E_{c/a}$ , and corresponding approximated martensitic phase transition temperature,  $T_M = \Delta E_{c/a}$  in K, nearest neighbor magnetic interaction constant,  $J_{ij}$ , with resulting mean-field Curie temperature  $T_C$ , and the energy difference between the ABCD and the ACBD order (see Fig. 1),  $\Delta E_{\text{order}}$ .

$c/a$	Fe <sub>2</sub> CoZn				Fe <sub>2</sub> CoGa		
	ABCD		ACBD		ABCD		ACBD
	0.92	1.4	1.0	1.36	0.88	1.46	1.0
$a$ (Å)	5.782		5.750		5.774		5.736
$M$ ( $\mu_B/\text{f.u.}$ )	6.65	6.64	5.95	6.24	6.07	6.08	5.29
$E_{\text{mix}}$ (meV/atom)	232	51	92	135	-257	-463	-642
$\Delta E_{c/a}$ (meV/atom)	44		-11		52		
$T_M$ (K)	483				563		
$J_{ij}$ (meV/atom)	23	20	19		18	14	23
$T_C$ (K)	1930	1725	1716		1540	1232	1645
$\Delta E_{\text{order}}$ (meV/atom)			-8.89				44.42

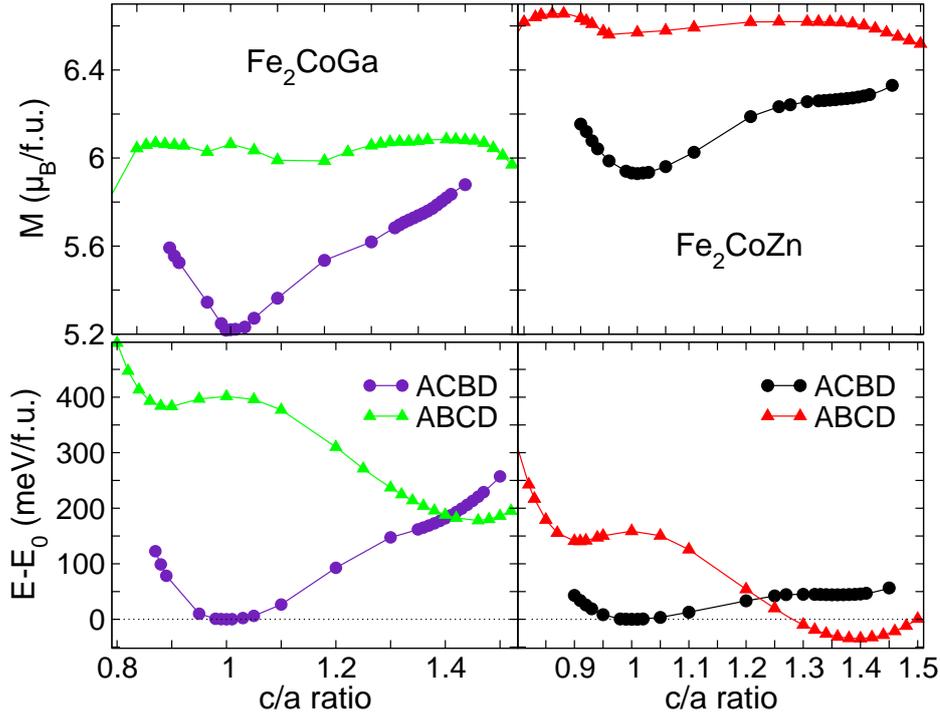
19×19×19 (mixing energies). The energy cut-off (in eV) was chosen as 366.5 for Fe<sub>2</sub>CoGa(Zn) and Co<sub>2</sub>FeZn.

The mixing energies are calculated using the simple formula:  $\Delta E_{\text{mix}} = E_{\text{alloy}} - \sum_i c_i E_i$ . Here  $E_{\text{alloy}}$  is the ground state energy of the respective alloy in eV/f.u. and  $E_i$  is the bulk energy in eV/atom of material component  $i$ , and  $c_i$  their respective weight in the formula unit. For the determination of the magnetic exchange interaction constants  $J_{ij}$ , we applied the Munich spin polarized relativistic Korringa-Kohn-Rostoker (SPR-KKR) package, version 3.6 [28,29]. The angular momentum expansion was truncated after the  $l = 3$  component. In the classical Heusler structure (order type ABCD, see Fig. 1) a k-point mesh of  $22 \times 22 \times 22$  and for the ACBD order a mesh of  $12 \times 12 \times 12$  was used. The Wigner-Seitz radii (in Å) are chosen for Fe<sub>2</sub>CoZn in ABCD order as: Fe = 2.669, Co = 2.628, and Zn = 2.787 and for ACBD order as: Fe = 2.654, Co = 2.613, and Zn = 2.772. For Fe<sub>2</sub>CoGa in ABCD order we get: Fe = 2.673, Co = 2.632, and Ga = 2.761 and in ACBD order we used Fe = 2.655, Co = 2.615, and Ga = 2.743.

### 3 Results and discussion

In order to investigate the structural stability of Fe-Co-Ga(Zn) alloys, we calculated the mixing energy,  $E_{\text{mix}}$ , and the energy difference between the classically Heusler order (type ABCD) and FeCoFeGa(Zn) order (type ACBD),  $\Delta E_{\text{order}}$ , both in meV/atom (compare Table 1 and Fig. 1). Thus, a negative  $\Delta E_{\text{order}}$  means that the ABCD order is energetically preferred. When performing a tetragonal  $c/a$  distortion of the ABCD structure, two energy minima are observed. The energy difference between these minima,  $\Delta E_{c/a}$ , expressed in K, gives a rough measure of the martensitic phase transition temperatures  $T_M$  (Table 1). For the ordering type ACBD with Co on half of the Fe sites (cf. Fig. 1), the cubic phase with  $c/a = 1$  is clearly favorable (see Fig. 2) From the calculated magnetic exchange interaction constants (Table 1 and Fig. 3) we evaluated the Curie temperatures using a simple Heisenberg model within mean field approximation:  $T_C = zM_i M_j J_{ij}$ , where  $z$  is the number of nearest neighbors and  $M_{i,j} = 1$  for ferromagnetic coupling. Here, we find extraordinarily high  $T_C$  values in the range of 1232 K up to 1645 K.

For Fe-Co-Ga the mixing energies are negative (cf. Table 1). And as early experimental studies confirm, these alloys are in fact stable [24]. The calculated  $J_{ij}$  predict strong ferromagnetism (cf. Fig. 3) and extremely high Curie temperatures of the order of  $1232 \text{ K} \leq T_C \leq 1645 \text{ K}$ . Also, for the case that Fe atoms become nearest neighbors as in order type ACBD, no antiferromagnetic tendency has been observed in the calculation. As is obvious from Fig. 2, the stable state at zero temperature is the cubic structure with order type ACBD in accordance with experiment.

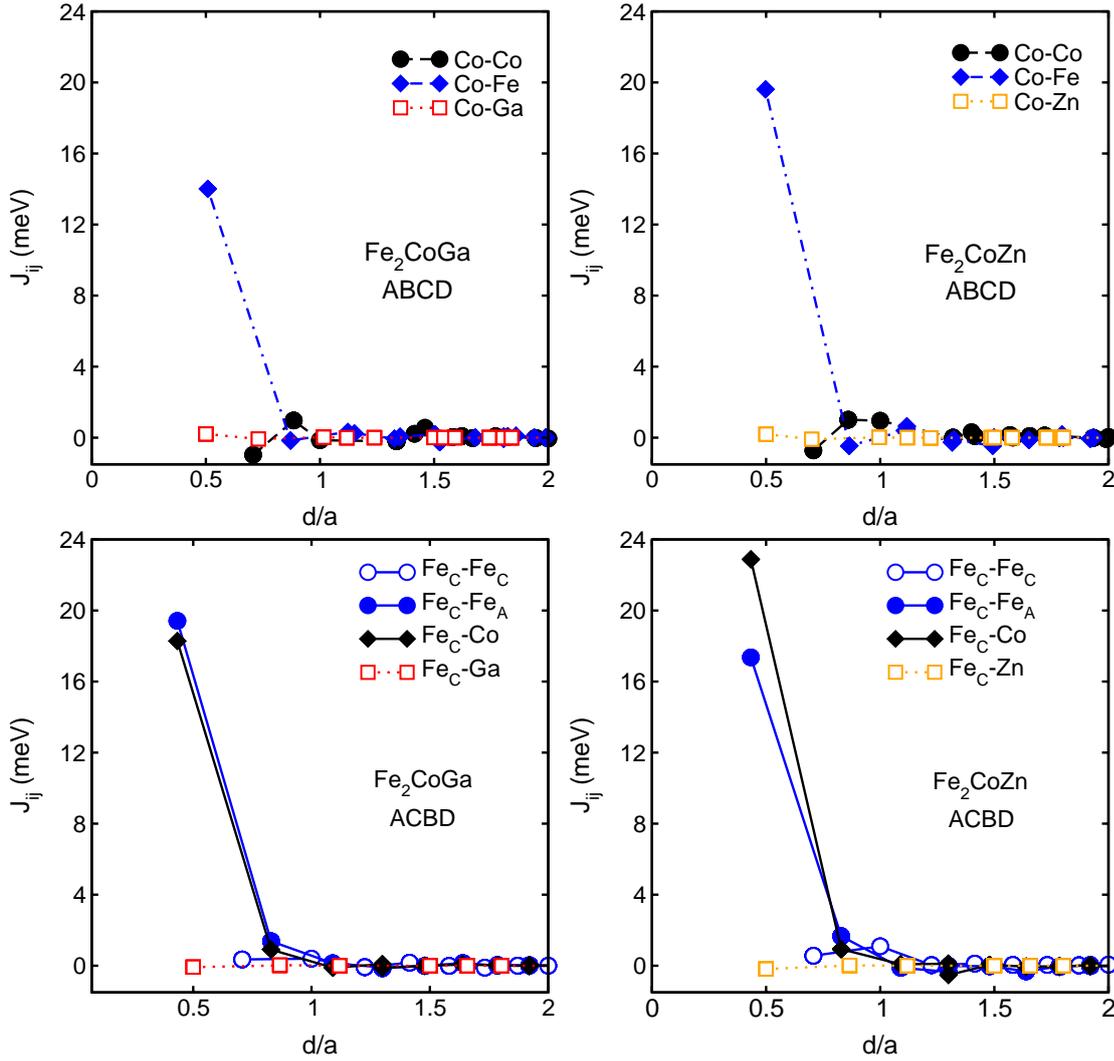


**Fig. 2.** The variation of the tetragonal distortion,  $c/a$ , and the corresponding magnetic moments,  $M$  (in  $\mu_B/\text{f.u.}$ ), for Fe-Co-Ga (left) and Fe-Co-Zn (right). Circles mark ordering type ACBD while triangles mark the classical Heusler in ordering type ABCD where Co and Ga(Zn) share one simple cubic sublattice (cf. Fig. 1). For ordered ABCD, two energy minima are found while for ACBD with Co on part of the Fe sublattice (B site), the cubic phase is clearly favoured.

Therefore, a martensitic phase transition is not likely to be expected in the Ga based alloy. Nevertheless assuming an artificial stabilization of the ABCD phase, we estimate a fairly high hypothetical  $T_M = 563$  K by expressing  $\Delta E_{c/a}$  in K. But at the same time, a phase transition would probably be accompanied with a diffusive rearrangement of Co atoms and thus be not martensitic. In order to confirm the proposed tendency of Co atoms occupying part of the Fe sites (e.g. B site), we performed analogous calculations for  $\text{Co}_2\text{FeZn}$  (not shown). And indeed, we again find a negative mixing energy but only one energy minimum at  $c/a = 1$ . The magnetic moment per formula unit is of the order  $6 \mu_B$  for the ABCD order and in between  $5.2$  and  $5.6 \mu_B$  for the ACBD structure.

For Fe-Co-Zn, the ABCD structure is found to be more favorable than the ACBD structure and our investigations would predict similar high Curie and martensitic temperatures (1725 K and 483 K, respectively, cf. Table 1). However, the calculated mixing energies appear to be positive for the classical Heusler structure (ABCD) as well as for the ACBD order. This means that at least for the stoichiometric composition, this system is not expected to be stable. But systematic alloying of Zn to Fe-Co-Ga may stabilize the ABCD structure and result in an interesting quaternary Fe-Co-Ga-Zn ferromagnetic shape memory alloy with high Curie and martensite temperatures. The magnetic moment per formula unit is enhanced by 8% in comparison to Fe-Co-Ga. This enhancement is due to an increased Fe moment of about 6% and an increased Co moment of about 11%. In summary, our result show that:

- In Fe-Co-Ga the mixing energies are negative for ABCD and ACBD order. An extraordinarily high  $T_C$  is predicted ( $\sim 1645$  K). But here the ACBD order is preferred over the ABCD order. Assuming a stabilisation of ABCD order a high  $T_M$  of 563 K can be expected. Still, not a martensitic transition but a diffusive rearrangement of Co atoms should be expected.
- The mixing energies for Fe-Co-Zn are positive. The magnetic moments are enhanced of about



**Fig. 3.** Magnetic exchange constants  $J_{ij}$  as a function of the distance  $d$  in units of the lattice constant  $a$  of  $\text{Fe-Co-Ga(Zn)}$ . The distance is measured from site C (cf. Fig. 1). Left upper panel: Exchange constants of  $\text{Fe}_2\text{CoGa}$  in ABCD order at  $c/a = 1.46$ . Left lower panel:  $\text{Fe}_2\text{CoGa}$  in ACBD order at  $c/a = 1$ . Right upper panel:  $\text{Fe}_2\text{CoZn}$  in ABCD order at  $c/a = 1.4$ . Right lower panel:  $\text{Fe}_2\text{CoZn}$  in ACBD order at  $c/a = 1$ . For the ABCD structures (upper panels) filled black circles denote the Co-Co interaction, filled blue diamonds mark the interaction between Co and Fe atoms, and open red (orange) squares belong to Co-Ga(Zn) interaction. For the ACBD order (lower panel), open blue circles mark the interaction between Fe atoms of the same type ( $\text{Fe}_C\text{-Fe}_C$ ) while filled blue circles denote the exchange interaction between nearest neighbor Fe atoms ( $\text{Fe}_C$  with  $\text{Fe}_A$ , see Fig. 1). Black filled diamonds belong to interaction of  $\text{Fe}_C$  with Co atoms while red (orange) open squares mark the interaction between  $\text{Fe}_C$  atoms and Ga (Zn) atoms. Strong ferromagnetism is obtained even if Fe atoms are nearest neighbors as for ACBD order.

8%. For a hypothetical Fe-Co-Zn Heusler alloy, the ABCD order with a tetragonal distorted  $c/a$  ratio of 1.4 would be more stable than the cubic ACBD structure. Strong ferromagnetism ( $T_C = 1725$  K) and high martensitic phase transition temperatures should be expected. But analogously to Fe-Co-Ga, a diffusive transition with a rearrangement of Co atoms may occur because the ACBD order with cubic symmetry lies only 8.8 meV/atom higher in energy than tetragonally distorted ABCD.

(c) For the stable  $\text{Co}_2\text{FeZn}$  system, the mixing energy is negative but only one energy minimum at  $c/a = 1$  is found.

These findings led us to the tentative conclusion, that, although Fe-Co-Zn is not stable, the addition of Ga to a hypothetical  $\text{Fe}_2\text{CoZn}$  Heusler alloy may lead to an interesting quaternary Fe-Co-Zn-Ga FSMA. But attention has to be paid to the possibility of a diffusive transition accompanied with a rearrangement of atoms, due to the fact that at least for stoichiometric  $\text{Fe}_2\text{CoGa}$ , Co atoms are expected to occupy part of the Fe sublattice (e.g., the B sites) and thus hinder the martensitic transformation. Further addition of Co leads again to negative mixing energies. But in the case of  $\text{Co}_2\text{FeZn}$ , the cubic system is the groundstate and no martensitic phase transition is expected. Thus, a careful variation of composition of the quaternary alloy  $\text{Fe}_{2-x}\text{Co}_{1+x}\text{Zn}_{1-y}\text{Ga}_y$  may result in a new ferromagnetic shape memory alloy with improved properties such as very high Curie and martensite temperature, eventually avoiding brittleness of prototype  $\text{Ni}_2\text{MnGa}$  Heusler alloy.

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### References

1. H.E. Karaca, I. Karaman, B. Basaran, Y.I. Chumlyakov, H.J. Maier, *Acta Materialia* **54**, 233 (2006)
2. M.A. Marioni, R.S. O'Handley, S.M. Allen, S.R. Hall, D.I. Paul, M.L. Richard, J. Feuchtwanger, B.W. Peterseon, J.M. Chambers, R. Techapiesancharoenkij, *J. Magn. Magn. Mater.* **290**, 35 (2005)
3. P. Müllner, V.A. Chernenko, G. Kostorz, *J. Magn. Magn. Mater.* **267**, 325 (2003)
4. R. Kainuma, F. Gejima, Y. Sutou, I. Ohnuma, K. Ishida, *Mater. Trans. JIM* **41**, 943 (2000)
5. Y. Sutou, Y. Imano, N. Koeda, T. Omori, R. Kainuma, K. Ishida, K. Oikawa, *Appl. Phys. Lett.* **85**, 4358 (2004)
6. W. Maziarz, J. Dutkiewicz, R. Santamarta, E. Cesari, *Eur. Phys. J. Special Topics* **158**, 137 (2008)
7. Y. Tanaka, K. Oikawa, Y. Sutou, T. Omori, R. Kainuma, K. Ishida, *Mater. Sci. Eng.* **438**, 1054 (2006)
8. P.J. Brown, K. Ishida, R. Kainuma, T. Kanomata, K.U. Neumann, K. Oikawa, B. Ouladdiaf, K.R.A. Ziebeck, *J. Phys.: Condens. Matter* **17**, 1301 (2005)
9. V.A. Chernenko, J. Pons, E. Cesari, I.K. Zaslachuk, *Scripta Materialia* **50**, 225 (2004)
10. K. Oikawa, T. Ots, F. Gejima, T. Ohmori, R. Kainuma, K. Ishida, *Mater. Trans.* **42**, 2472 (2001)
11. K. Oikawa, L. Wulff, T. Iijima, F. Gelima, T. Ohmori, A. Fujita, K. Fukamichi, R. Kainuma, K. Ishida, *Appl. Phys. Lett.* **79**, 3290 (2001)
12. M. Wuttig, J. Li, C. Craciunescu, *Scripta Materialia* **44**, 2393 (2002)
13. C. Craciunescu, Y. Kishi, A.T. Lograsso, M. Wuttig, *Scripta Materialia* **47**, 285 (2002)
14. K. Oikawa, T. Ota, Y. Tanaka, T. Omori, R. Kainuma, K. Ishida, *Trans. Mater. Res. Soc. Jpn.* **28**, 265 (2003)
15. N. Koeda, Y. Sutou, T. Omori, K. Oikawa, R. Kainuma, K. Ishida, *Scripta Materialia* **52**, 1153 (2005)
16. K. Oikawa, N. Koeda, Y. Sutou, T. Omori, R. Kainuma, K. Ishida, *Mater. Trans.* **45**, 2780 (2004)
17. H. Morito, A. Fujita, K. Oikawa, K. Fukamichi, R. Kainuma, T. Kanomata, K. Ishida, *J. Phys.: Condens. Matter* **21**, 076001 (2009)
18. Y.I.T. Omori, K. Oikawa, Y. Sutou, R. Kainuma, K. Ishida, *Mater. Sci. Eng.* **438**, 970 (2006)
19. T. Takayama, S. Shinohara, K. Ishida, T. Nishizawa, *J. Phase Equilibria* **16**, 390 (1995)
20. J.V. Humbeeck, J. Janssen, N. Mwamba, L. Delaey, *Scripta Metall.* **18**, 893 (1984)
21. J. Gil, J.M. Guilemany, *J. Mater. Sci. Lett.* **11**, 493 (1992)
22. H. Numakura, M. Ishimoto, M. Wuttig, *Mater. Sci. Eng. A* **442**, 195 (2006)
23. Z. Lei, J. Cheng-Bao, S. Jia-Xiang, X. Hui-Bin, *Chinese Physics B* **18**, 1647 (2009)
24. N.K. Jaggi, K.R.P.M. Rao, A.K. Grover, L.C. Gupta, R. Vijayaraghavan, L.D. Khoi, *Hyperfine Interactions* **4**, 402 (1978)
25. G. Kresse, J. Furthmüller, *Phys. Rev. B* **54**, 11169 (1996)

26. G. Kresse, D. Joubert, Phys. Rev. B **59**, 1758 (1999)
27. J.P. Perdew, K. Burke, M. Ernzerhof, Phys. Rev. Lett. **77**, 3865 (1996)
28. H. Ebert, in *Lecture Notes in Physics*, edited by Springer (H. Dreyse, Berlin, 2000), Vol. 50, p. 191
29. Ebert et al, <http://olymp.cup.uni-muenchen.de/ak/ebert/SPRKKR>